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# Dynamically modelling the MW

# Outline

- The challenge
- Equilibrium dynamical models
- Axisymmetric models
- Including the bar
  - M<sub>2</sub>M models
  - Torus models
- Why orbit-based models of the solar nhd are problematic
- Conclusions

# The challenge

- Spectacular data sets are flooding in
  - Gaia + spectroscopic surveys
    - APOGEE, RAVE, GES, LAMOST, Galah, WEAVE,....
- We want to learn how our Galaxy formed
  - First we must establish how it is structured now
- Each survey has non-trivial selection function
  - It images the MW through a distorting lens
- Also most interesting parts of the MW obscured by dust
- So non-trivial to infer MW's structure from catalogues
- We must build a model that's consistent with catalogues
  - This model will *embody all we know about the MW*

# Models

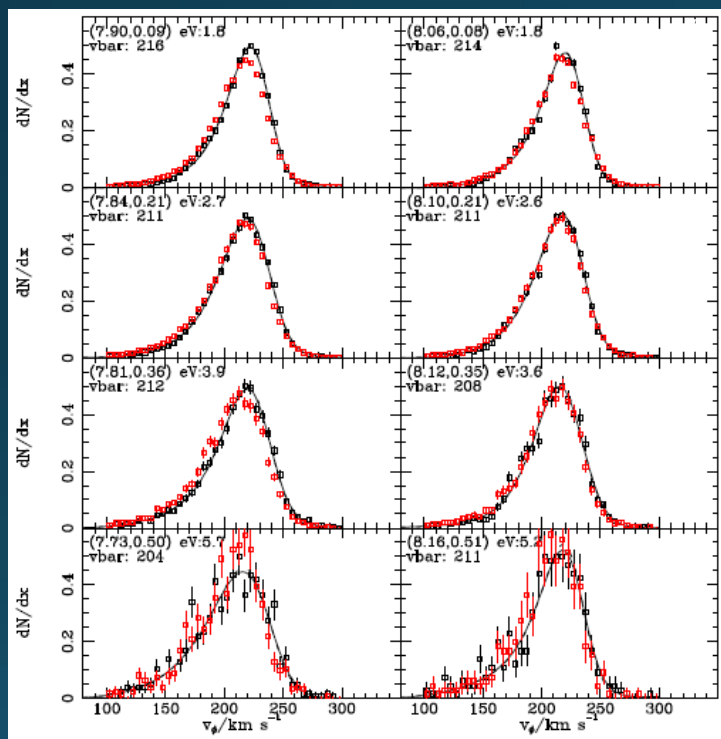
- We have to model stars & gas in parallel
  - Don't forget: since 1950s radio observations of HI, CO, etc have provided crucial constraints on MW's gravitational field  $g(x)$
- >90% of MW's mass is dark
  - We must model DM too!
- We have to track DM from its contribution to  $g(x)$ 
  - Can only infer  $g(x)$  from kinematics to the extent that the MW is in equilibrium
- So we must start from *equilibrium dynamical* models
  - Later we will add non-equilibrium features (spirals, warp, streams,..) as perturbations

# Jeans theorem

- Given equilibrium, Jeans' theorem collapses 6d phase space to 3d orbit space
- In principle infinite freedom in choice of constants of motion to use as coordinates of orbit space
  - But in practice only one rational choice:
    - action integrals  $J$
- We now have axisymmetric models specified by analytic
  - Extended Distribution Functions (EDFs)  $f_*(J, \text{age}, [\text{Fe}/\text{H}])$
  - $f_{\text{DM}}(J)$
  - (B & McMillan 2011; Bovy & Rix 2013; Sanders & B 2015; Penoyre+ 2015)

# Rave kinematics Binney, Burnett + RAVE 2014

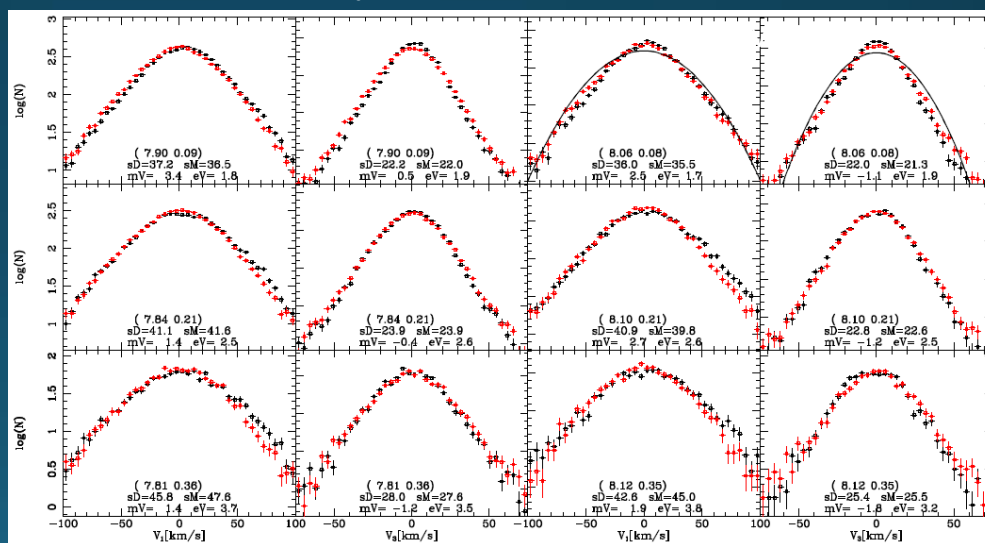
- Binney (2012) fitted disc  $f(J)$  to GCS data ( $s \llsim 0.1$  kpc)
- Binney + (2014) tested its *predictions* for kinematics of RAVE stars in 8 volumes with  $s \llsim 2$  kpc



## Cool dwarfs

$R < R_0$

$R > R_0$



# Self-consistent multi-component models

- First models used fixed parametrised  $\Phi$  but current models have self-consistent  $\Phi$
- Major advantages of using actions:
  - Ease of constructing models that contain many populations (stars of each [mass, age & chemistry], DM, white dwarfs, n-stars,...)
  - Ease of finding self-consistent gravitational potential  $\Phi(x)$ 
    - Penoyre+ 2015; B & Piffl 2015; Cole & B 2016

# Non-axisymmetry

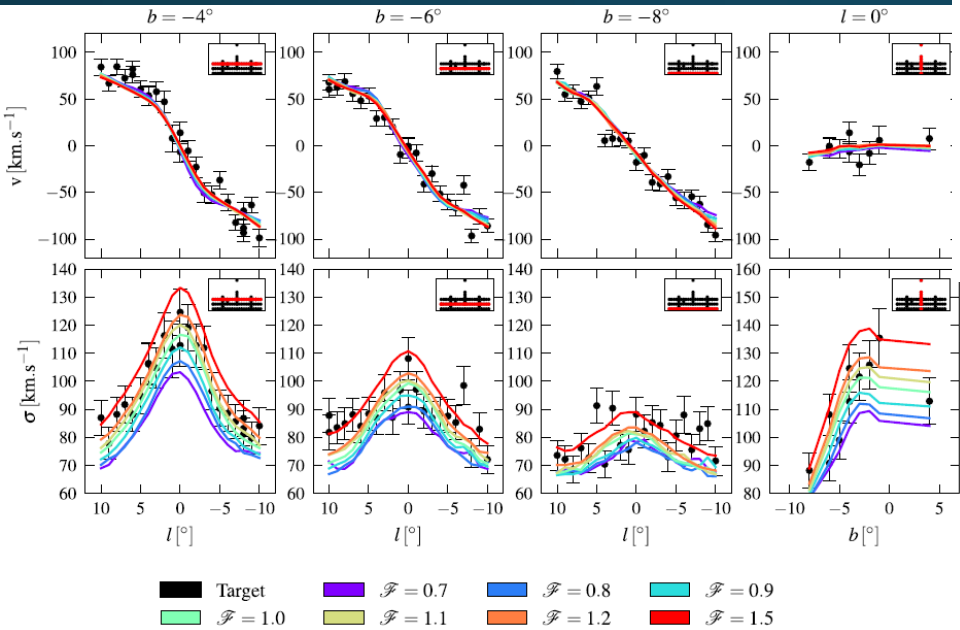
- MW is a *barred spiral*
  - Simplest ansatz is that  $\Phi$  steady in rotating frame
- Models to date exploit Staeckel Fudge, which gives  $J(x,v)$ 
  - Converts  $f(J)$  to  $f(x,v)$  so can get  $v$ -distribution &  $\rho$  by integrating over  $v$
- Unfortunately SF not available for  $\Phi$  with figure rotation
- Without SF can only build orbit-based models
  - N-body, M2M, Schwarzschild, Torus
- We treat a galaxy as a sum of orbits



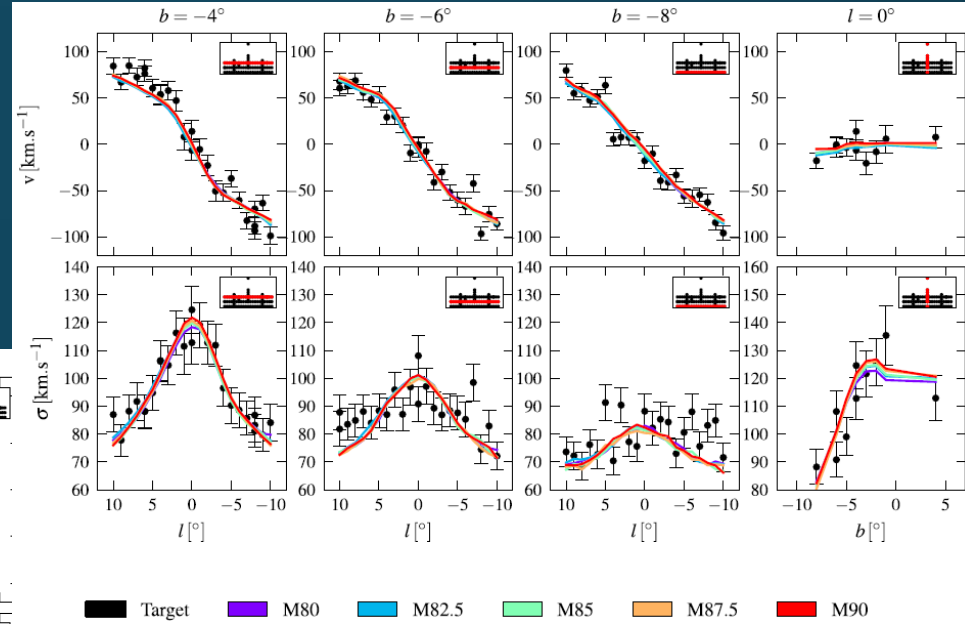
# M<sub>2</sub>M models

- (Syer & Tremaine 1996; De Lorenzi+ 2007; Long+ 2013; Martinez-Valpuesta 2012; Portail+ 2015)
- Let N-body model develop a bar
- Vary particle weights while integrating eqs of motion to optimize fits to observables
- Occasionally update  $\Phi$
- Has produced very convincing models of bar/bulge

Pattern speed: 27 km/s/kpc



Increasing stellar mass



Decreasing halo mass in bulge

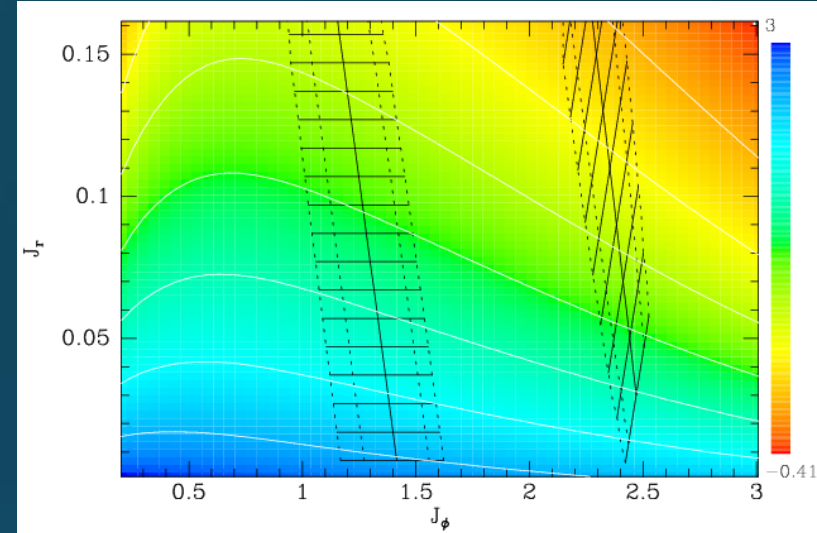
Portail+ 2015

# Schwarzschild & Torus modelling

- Schwarzschild (1979) is industry standard for models of early-type galaxies
- Choose  $\Phi$  and integrate an orbit library
  - then choose weights to fit data
- For many reasons it's best to represent orbits with tori
  - Structures with known actions  $J$  equipped with angle variables
  - A torus encodes every orbit  $J$  to all eternity in  $\sim 100$  numbers
  - Torus Mapper (B & McMillan 2016) generates torus more quickly than a long orbit integration
  - TM's methods return anything you might compute ( $[\rho(x), v(x), \dots]$ )
  - Can now relax  $\Phi$  to self-consistency
- Traditional representation of orbits as time series not competitive

# Orbit families

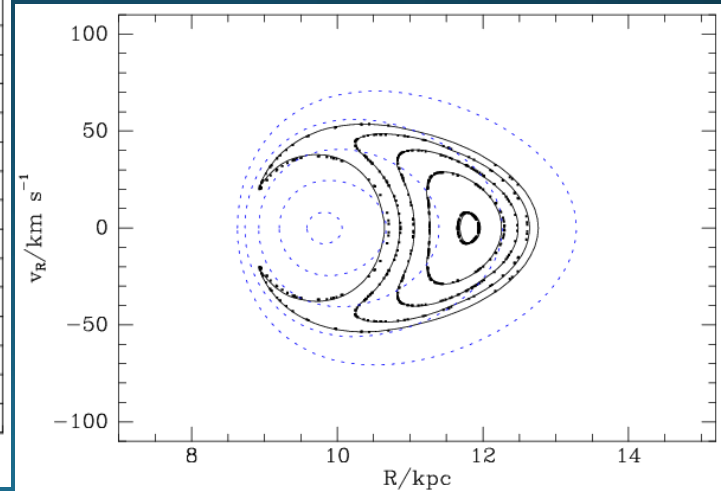
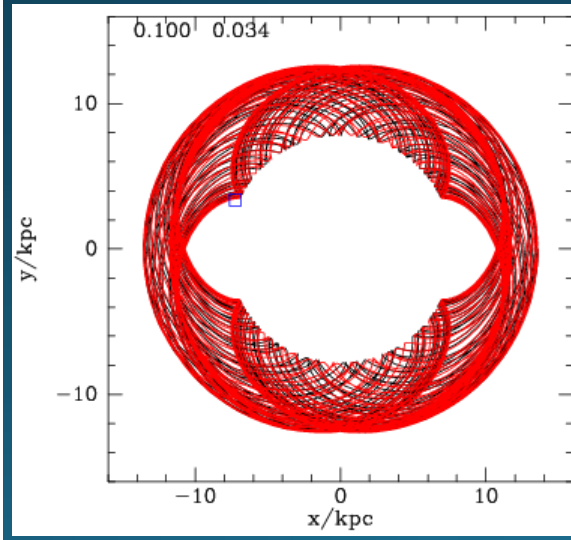
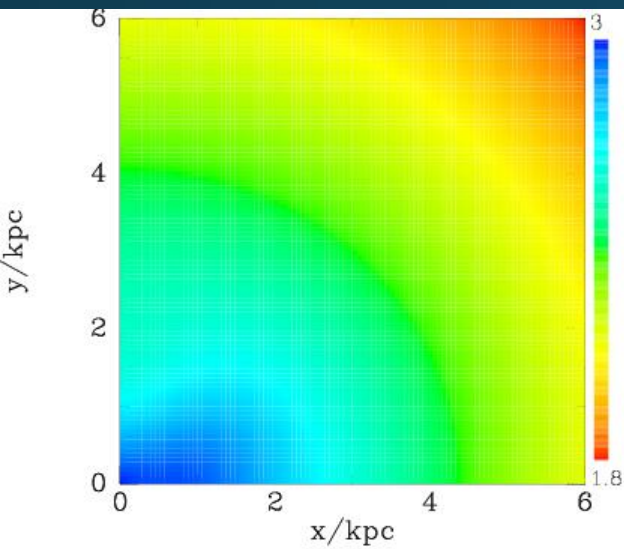
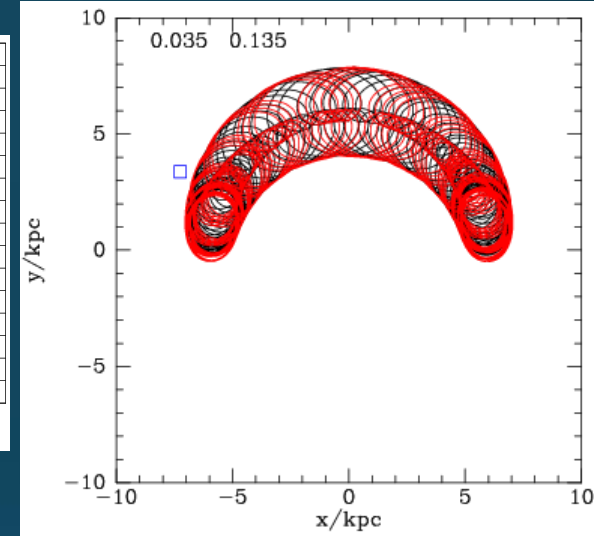
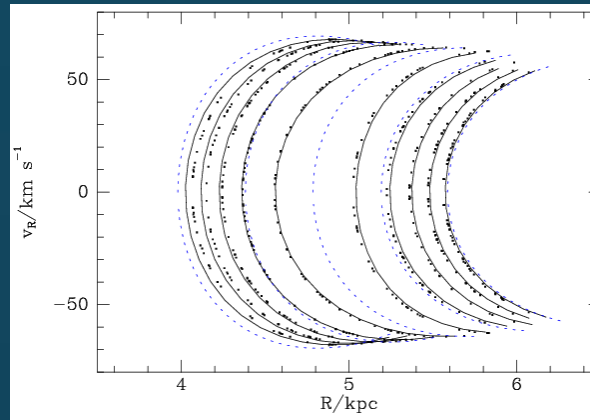
- Once  $\Phi$  is non-axisymmetric, cannot ignore that there's  $>1$  orbit family
- Each family has its own angle-action coordinates
- Orbits that belong to  $>1$  family are "chaotic"
- Families normally considered to arise from "resonant trapping"
- In a realistic Galactic bar (Sormani+2016) corotation and Lindblad resonances have significant zones of entrapment
- Using perturbation theory we can construct tori for trapping zones



Binney 2017

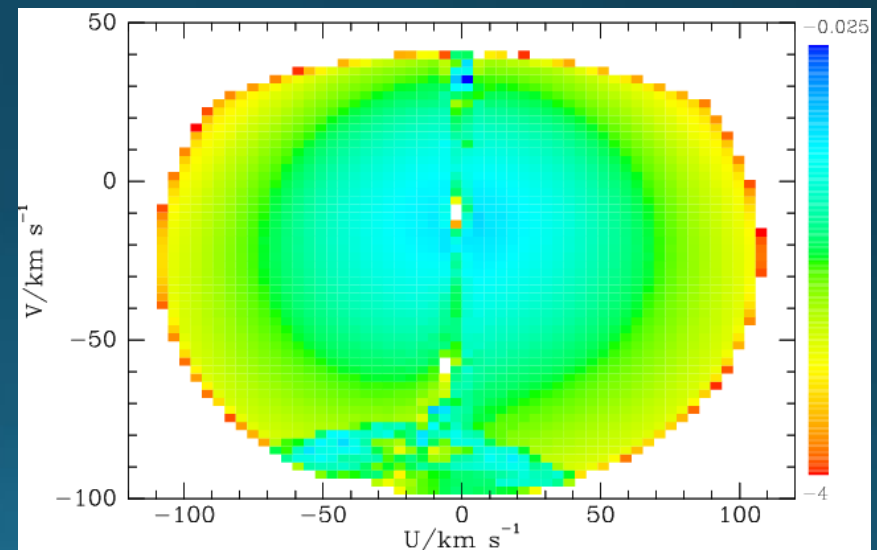
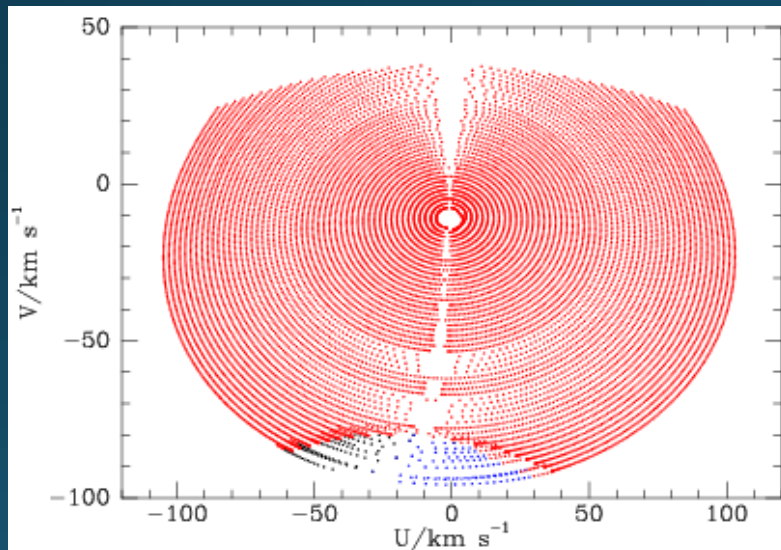
# Non-axisymmetric tori

(Binney 2017)



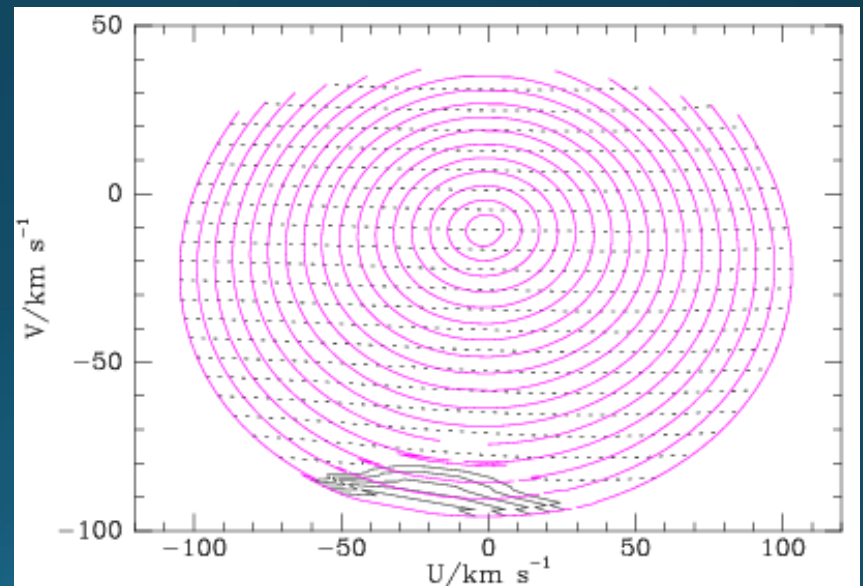
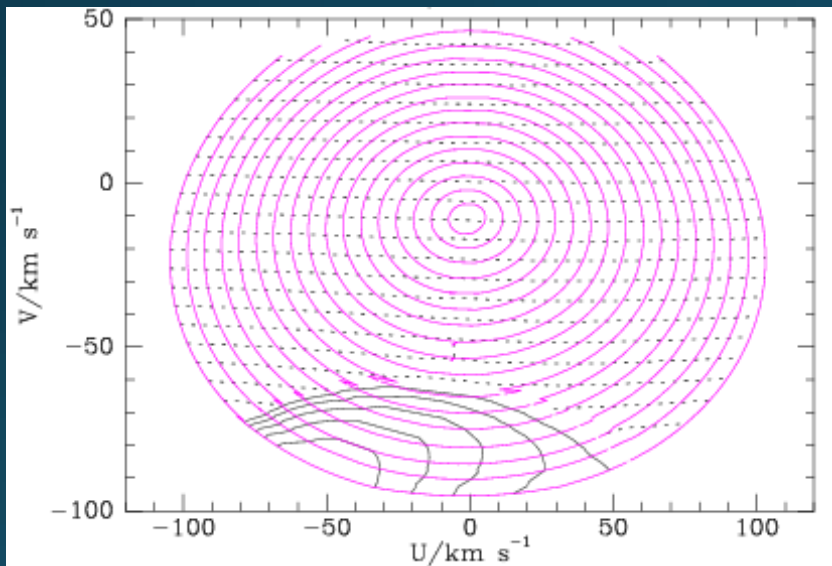
# Local v-space from orbits

- Individual tori are beautiful
- But still hard to recover kinematics of solar nhd
  - Uniform grid in action space maps to an irregular grid in velocity space
  - Density of an orbit diverges at its edges
  - So small # of orbits contribute heavily to certain Vs -> enhanced Poisson noise



# New constants of motion

- Determine map  $v \rightarrow J$
- Can then evaluate  $f(v)$  (Schoenrich & McMillan 2016)



# Conclusions

- Only by building models can we understand how the MW is structured
- We should start from equilibrium dynamical models
- Then Jeans theorem puts MW into 3d orbit space
- Actions are the preferred coordinates for orbit space
- DFs  $f(J)$  for DM and various stellar types readily combined to build multi-component, self-consistent axisymmetric models
- The bar forces recognition of several orbit families, each with its own AA variables
- Reluctantly accept that barred models cannot be based on  $J(x,v)$
- Excellent models of bar/bulge constructed by M2M technique
- Torus modelling supersedes Schwarzschild modelling as an alternative
- But any orbit model liable to excessive Poisson noise
- Use tori to determine map  $v \rightarrow J$  at various  $x$